

BASICS OF CONTINUUM MODELS

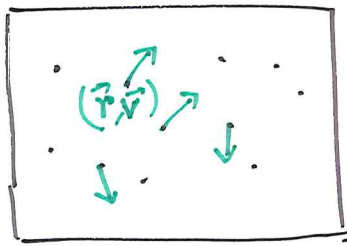
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Plasma is a complex medium consisting of electrons, ions and neutrals which moves under the influence of electric, magnetic and collisional forces. To understand plasma and its basis of dynamics, we'll look into classical continuum mechanics for neutral fluid and gases.

For single particles: VECTOR QUANTITIES ($\vec{r}, \vec{v}, \vec{F}$)

For Continuum models: VECTOR AND SCALAR FIELDS

Instead of tracking individual particles, we study fields which represent an ensemble of particles.

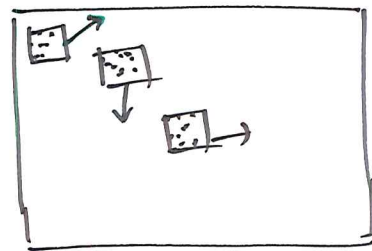


SINGLE

$$(\vec{r}, \vec{v})$$

↓

$$\vec{v} \text{ is not } f(\vec{r})$$



CONTINUUM/FLUID

$$(\vec{v}, \vec{E}, \vec{B}, \dots)$$

$$(\rho, T, \dots)$$

↓

$$\vec{v}(\vec{r}, t)$$

Velocity field is a function of space.

2

For continuum model, we'll consider density (ρ) instead of mass (m) unlike the single particle motion.

CHANGE IN DENSITY [$\rho(\vec{r}, t)$]

$$\underbrace{\frac{\partial}{\partial t} \rho}_{\text{Time}} \quad \underbrace{\vec{\nabla} \rho = \left(\frac{\partial}{\partial x} \hat{x} + \frac{\partial}{\partial y} \hat{y} + \frac{\partial}{\partial z} \hat{z} \right) \rho}_{\text{SPACE}}$$

RECALLING GAUSS THEOREM (DIVERGENCE THEOREM)

$$\boxed{\iiint_V \vec{\nabla} \cdot \vec{u} \, dV = \iint_S \vec{u} \cdot \hat{n} \, dS}$$

\downarrow
Field vector
 \downarrow
Surface normal

GAUSS LAW IN ELECTROSTATICS

$$\boxed{\int_S \vec{E} \cdot d\vec{s} = \frac{Q}{\epsilon_0}}$$

$$\rightarrow 4\pi r^2 \hat{r} E(r) = \frac{Q}{\epsilon_0}$$

$$\boxed{\Rightarrow E(r) = \frac{Q}{4\pi\epsilon_0 r^2}}$$

using Divergence theorem

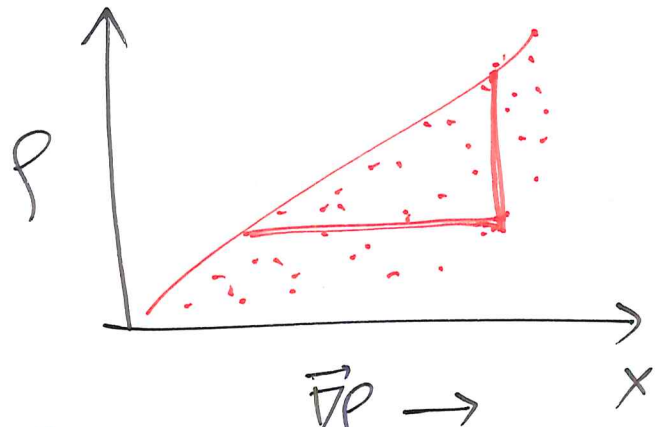
$$\Rightarrow \vec{\nabla} \cdot \vec{E} = \frac{\rho_e}{\epsilon_0}$$

$$\nabla^2 \phi = - \frac{\rho_e}{\epsilon_0}$$

where, $E = -\nabla\phi$

DENSITY FIELD (GENERAL FLUID)

IN 1D:



DISCRETE APPROACH

$\nabla \rho \rightarrow$
 Δ is large.

1) Density variation:

$$\rho(x + \Delta x) - \rho(x) = \Delta \rho$$

$$\Delta \rho > 0 \quad \text{or} \quad \Delta \rho < 0$$

For a stationary case where we only have spatial gradient.

2) Dynamic density (density moving)

$$U = \frac{\Delta x}{\Delta t}$$

\swarrow
 Velocity field
 (Fluid flow)

\searrow
 density element

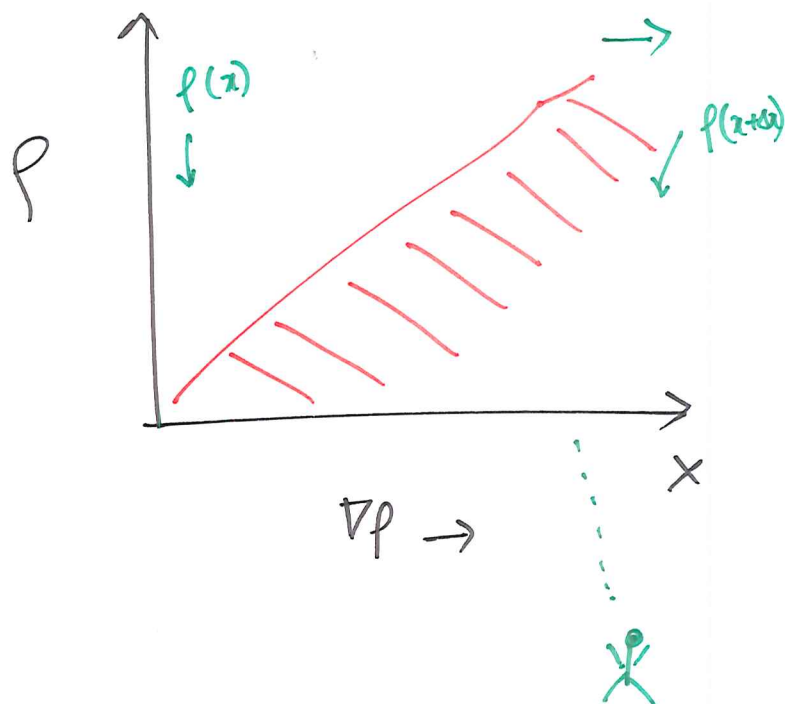
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FIXED POINT OBSERVER
OR

LABORATORY FRAME OF REFERENCE
OR

OR

STANDING OBSERVER



The observer will see
 $-\Delta p$ in time

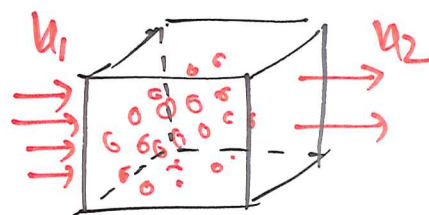
$$\frac{\Delta p}{\Delta t} = -u \frac{\Delta p}{\Delta x}$$

$$\Rightarrow \frac{\partial p}{\partial t} = \dot{p} = -u \frac{\partial p}{\partial x}$$

FOR SITUATIONS WHEN WE HAVE CHANGE IN VELOCITY

i.e. $\vec{\nabla} u \neq 0$

means more particle packed
together over time that
leads to compressibility.



$u_1 > u_2$

if more incoming flow
than outgoing, there will
be increase in density over time.

NOTE: For compressibility, we need not to have ∇f . Only $\nabla u \neq 0$ should be enough.

$$u(x + \Delta x) - u(x) = \Delta u$$

For a fixed observer, $\Delta u < 0$

$$\boxed{\frac{\Delta \rho}{\Delta t} = -\rho \frac{\Delta u}{\Delta x}} \quad \text{CONSERVATION OF MASS}$$

REMEMBER! For a fixed observer, (incoming field - outgoing field) to decide the sign.

CONSERVATION OF MASS IN 1D

$$M = \rho \Delta x \quad \Delta(\rho \Delta x) = (\rho u_{\text{in}} - \rho u_{\text{out}}) \Delta t$$

$$\Rightarrow \boxed{\frac{\Delta \rho}{\Delta t} = -\rho \frac{\Delta u}{\Delta x}}$$

IN GENERAL,

$$\boxed{\frac{\partial \rho}{\partial t} = -\frac{\partial}{\partial x} (\rho u)}$$

~~1D/3D/1D~~

IN 3D:

CONTINUITY EQN:

$$\frac{\partial}{\partial t} \rho + \vec{\nabla} \cdot (\vec{u} \rho) = 0$$

mass density.

mass flux

When there is no production or loss

Change in mass density is equal to: total change in mass flux in space.

Divergence of mass flux.

TAKING \int volume integral \Rightarrow mass conservation.

MOVING OBSERVER:

$$\frac{D}{Dt} = \frac{\partial}{\partial t} + \vec{u} \cdot \vec{\nabla}$$

Convective derivative

The variation in the flow due to convection.

$$\frac{D}{dt} \rho = -\rho \vec{\nabla} \cdot \vec{u}$$

compressibility

Change in time + movement of the system

IF FLUID IS INCOMPRESSIBLE

$$\frac{D}{Dt} \rho = 0$$

OLD \rightarrow $\left[\frac{\partial}{\partial t} \rho + \vec{u} \cdot \frac{\partial \rho}{\partial \vec{x}} = 0 \right]$

FOR A COMOVING OBSERVER NO DENSITY GRADIENT.

IN A COMOVING FRAME, ANY CHANGES ARE DUE TO EXPANSION OR COMPRESSION OF THE FLUID.

\downarrow
 $-\rho \nabla \cdot \vec{u}$

IF PRODUCTION OR LOSS

$$\frac{\partial}{\partial t} \rho + \nabla \cdot (\vec{u} \rho) = \alpha - \beta$$

\downarrow production \downarrow loss

For Plasma, (charge continuity)

$$\frac{\partial}{\partial t} (q \cdot n) + \nabla \cdot \vec{j} = 0 \quad \vec{j} = q n \vec{u}$$

- FIXED POINT OBSERVER \rightarrow EULERIAN APPROACH $\downarrow \frac{\partial}{\partial t}$
- COMOVING OBSERVER \rightarrow LAGRANGIAN APPROACH $\downarrow \frac{D}{Dt}$

VELOCITY FIELD

8

NEWTON'S SECOND LAW →

FOR SINGLE PARTICLE

$$m\vec{a} = \vec{F} = m \frac{d\vec{v}}{dt} \rightarrow \text{velocity of particle.}$$

IN FLUID,

$$\vec{u} = \vec{u}(\vec{r}, t)$$

position → time →

$\vec{v}(t)$
Function of time only

IN FLUID,

velocity → velocity field
mass → mass density field

$$\rho \frac{\partial \vec{u}}{\partial t} = \vec{f}$$

Force density

$$\vec{f} = \frac{\vec{F}}{V}$$

Force per volume

This is valid if \vec{u} is const./homogeneous at $t=0$

LAGRANGIAN.

$$\rho \frac{D}{Dt} \cdot \vec{u} = \rho \left(\frac{\partial}{\partial t} \vec{u} + \vec{u} \cdot \nabla \vec{u} \right) = \vec{f}$$

CONSIDERING CHANGE IN MOMENTUM IN A SMALL VOLUME

$$\frac{d}{dt} \int_V \rho \vec{u} d\vec{r} = - \oint_S \rho \vec{u} \vec{u} \cdot \hat{n} dS + \int_V \vec{f} d\vec{r}$$

using GAUSS THEOREM

9

$$= \int_V \left[\frac{\partial}{\partial t} (\rho \vec{u}) + \nabla \cdot \rho \vec{u} \vec{u} \right] d\vec{r}$$

$$= \int_V \left[\underbrace{\vec{u} \frac{\partial}{\partial t} \rho + \rho \frac{\partial}{\partial t} \vec{u}}_{\text{CONTINUITY EQN}} + \underbrace{\nabla (\rho \vec{u}) \vec{u} + (\rho \vec{u}) \cdot (\nabla \vec{u})}_{\text{For no production or loss} \rightarrow 0} \right] d\vec{r}$$

CONTINUITY EQN \rightarrow For no production or loss $\rightarrow 0$

$$= \int_V \left[\rho \frac{\partial \vec{u}}{\partial t} + (\rho \vec{u}) \cdot \nabla \vec{u} \right] d\vec{r}$$

$$\Rightarrow \int_V \rho \left(\frac{\partial}{\partial t} \vec{u} + \vec{u} \cdot \nabla \vec{u} \right) d\vec{r} = \int_V \vec{f} d\vec{r}$$

CONTINUITY
MOMENTUM

At this point, we have two eqns and three unknowns (ρ, \vec{u}, \vec{f})

Force \vec{f} can be represented as

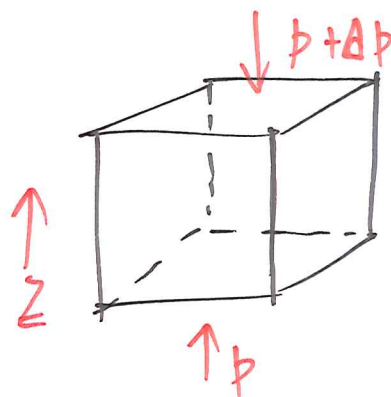
(10)

$$\vec{f} = -\vec{\nabla} p$$

Force
↓

$$F = A \cdot (p(z) - p(z+\Delta z))$$

$$\approx -A \cdot \Delta z \frac{\Delta p}{\Delta z}$$



$$\frac{F}{A \cdot \Delta z} = -\frac{\Delta p}{\Delta z} = -\vec{\nabla} p$$

$$\Rightarrow \vec{f} = -\vec{\nabla} p$$

THE MOST RELEVANT FORM OF MOMENTUM EQN FOR FLUID

$$\rho \left(\frac{\partial}{\partial t} \vec{u} + \vec{u} \cdot \vec{\nabla} \vec{u} \right) = -\underbrace{\nabla p}_{\text{pressure}} + \underbrace{\mu \nabla^2 \vec{u}}_{\text{viscosity}} + f$$

NEXT \rightarrow EQN OF STATE